

COL7160 : Quantum Computing

Lecture 18 : Grover's Search Algorithm

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1 Introduction

This lecture introduces Grover's Search Algorithm [Gro96]. It is the foundation for a family of related algorithms sharing the same core idea: Amplitude Amplification, Amplitude Estimation, Approximate Counting, and Quantum Random Walks. Unlike Simon's algorithm, which exploits algebraic structure in f , Grover's algorithm makes no structural assumption on f and still provides a provable quantum speedup.

2 The Search Problem

Definition 1 (Search Problem). Given a string $x \in \{0, 1\}^N$ and oracle access to $f : \{0, 1\}^n \rightarrow \{0, 1\}$ with $N = 2^n$, find x^* such that $f(x^*) = 1$. There are $t \geq 1$ such inputs.

In the *unique search* case, $t = 1$:

$$\exists i \text{ s.t. } x_i = 1, \quad \forall j \neq i, x_j = 0.$$

Classical complexity. A deterministic algorithm requires $\Omega(N)$ queries. A randomised algorithm requires $\Omega(N/t)$ queries for constant success probability. Grover's algorithm achieves $O(\sqrt{N/t})$ quantum queries.

Remark 2. The circuit for f is known but is useless for finding x^* directly; the only use is to construct the oracle. A canonical example is SAT: evaluating f on any assignment is efficient, but finding a satisfying assignment is computationally hard.

2.1 Quantum RAM

To create the superposition $\frac{1}{\sqrt{N}} \sum_i |i\rangle |x_i\rangle$ when x_i cannot be expressed as an efficient function of i , the data must be stored in **Quantum Random Access Memory (QRAM)**:

$$\frac{1}{\sqrt{N}} \sum_{i=0}^{N-1} |i\rangle |0\rangle \xrightarrow{\text{QRAM}} \frac{1}{\sqrt{N}} \sum_{i=0}^{N-1} |i\rangle |x_i\rangle.$$

For most applications, $x_i = f(i)$ for an efficiently computable f , so QRAM is not needed.

3 Grover's Algorithm

3.1 Setup and Key States

Prepare the uniform superposition:

$$|v\rangle = H^{\otimes n} |0\rangle^{\otimes n} = \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle.$$

Definition 3 (Good and Bad States).

$$|A\rangle = \frac{1}{\sqrt{t}} \sum_{\substack{x \\ f(x)=1}} |x\rangle, \quad |B\rangle = \frac{1}{\sqrt{N-t}} \sum_{\substack{x \\ f(x)=0}} |x\rangle.$$

$|A\rangle$ is the uniform superposition over marked (good) states; $|B\rangle$ over unmarked (bad) states. They are orthonormal.

The initial state decomposes as:

$$|v\rangle = \sqrt{\frac{t}{N}} |A\rangle + \sqrt{\frac{N-t}{N}} |B\rangle = \sin \theta |A\rangle + \cos \theta |B\rangle, \quad \sin \theta = \sqrt{\frac{t}{N}}.$$

3.2 The Phase Oracle

$$Z_f |x\rangle = (-1)^{f(x)} |x\rangle.$$

Implemented via the bit-flip oracle $\mathcal{O}_f : |x\rangle |b\rangle \mapsto |x\rangle |b \oplus f(x)\rangle$ with ancilla in $|-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$ (phase kickback). Circuit size equals that of f .

3.3 The Diffusion Operator: Flip Around the Mean

Definition 4 (Diffusion Operator).

$$D = 2|v\rangle\langle v| - I = H^{\otimes n}(2|0\rangle\langle 0| - I)H^{\otimes n}.$$

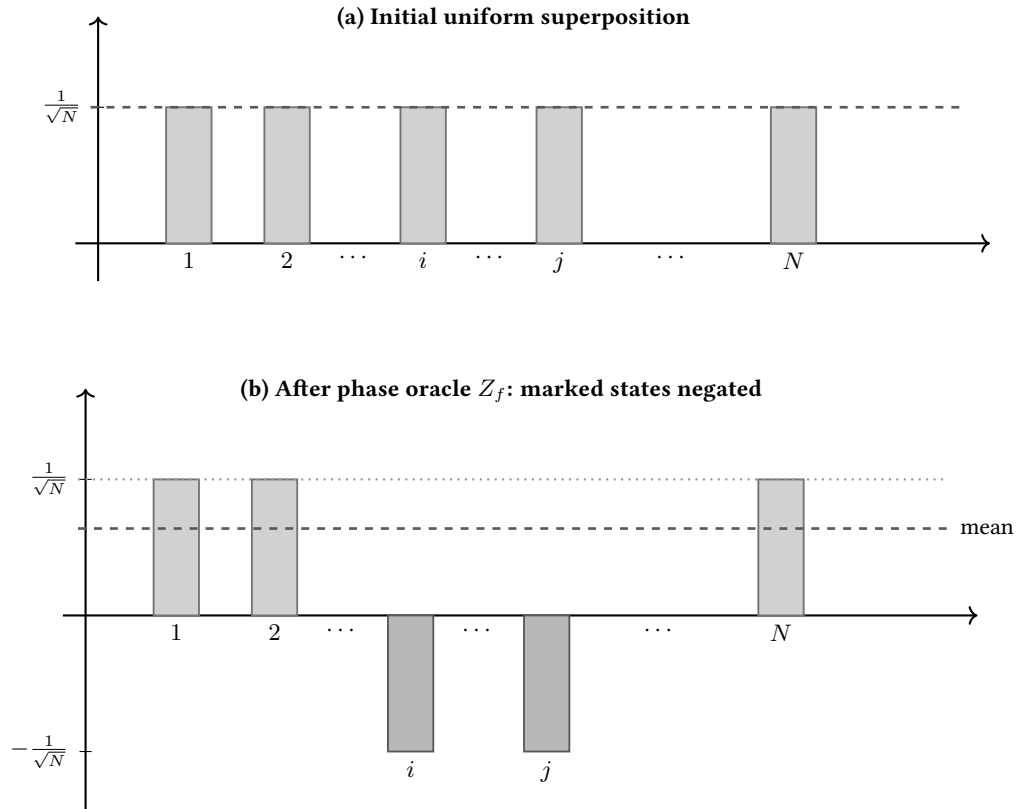
Lemma 5. If $|\psi\rangle = \sum_x \alpha_x |x\rangle$ with mean $\bar{\alpha} = \frac{1}{N} \sum_x \alpha_x$, then $D|\psi\rangle = \sum_x (2\bar{\alpha} - \alpha_x) |x\rangle$.

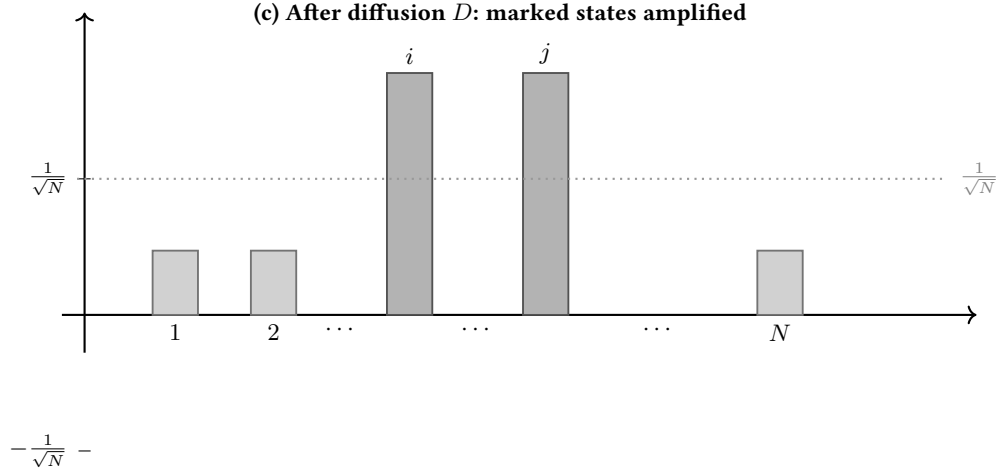
Proof. Since $\langle v|\psi\rangle = \sqrt{N} \bar{\alpha}$,

$$D|\psi\rangle = 2|v\rangle\langle v|\psi\rangle - |\psi\rangle = \sum_x (2\bar{\alpha} - \alpha_x) |x\rangle. \quad \square$$

Thus D reflects each amplitude about the mean: $\alpha_x \mapsto 2\bar{\alpha} - \alpha_x$.

The three figures below illustrate one complete Grover iteration. Indices i and j are the marked states ($f(x) = 1$); all others are unmarked.





In (a), all amplitudes equal $1/\sqrt{N}$ and the mean equals $1/\sqrt{N}$. In (b), the phase oracle negates the amplitudes of the marked states i and j to $-1/\sqrt{N}$, causing the mean to drop below $1/\sqrt{N}$. In (c), the diffusion operator maps $\alpha_x \mapsto 2\bar{\alpha} - \alpha_x$: the marked states (at $-1/\sqrt{N}$, far below the mean) are reflected to a value well above $1/\sqrt{N}$, while unmarked states (slightly above the mean) decrease. Repeated application increases the probability of measuring a marked state.

3.4 The Grover Operator

Definition 6 (Grover Operator).

$$G = D \cdot Z_f = (2|v\rangle\langle v| - I) \cdot Z_f.$$

Applying Z_f to $|v\rangle$ gives:

$$Z_f |v\rangle = -\sqrt{\frac{t}{N}} |A\rangle + \sqrt{\frac{N-t}{N}} |B\rangle.$$

3.5 Geometric Analysis

Theorem 7. *In the plane $\text{span}\{|A\rangle, |B\rangle\}$, each application of G rotates the state vector by 2θ towards $|A\rangle$. After k applications:*

$$G^k |v\rangle = \sin((2k+1)\theta) |A\rangle + \cos((2k+1)\theta) |B\rangle.$$

Proof. The initial state $|v\rangle = \sin \theta |A\rangle + \cos \theta |B\rangle$ lies at angle θ from $|B\rangle$.

Step 1 – Z_f : The phase oracle negates the $|A\rangle$ component, acting as the reflection $I - 2|A\rangle\langle A|$ in this plane. This maps the state from angle θ to $-\theta$ (reflection about $|B\rangle$):

$$Z_f |v\rangle = -\sin \theta |A\rangle + \cos \theta |B\rangle.$$

Step 2 – D : The diffusion $D = 2|v\rangle\langle v| - I$ reflects about $|v\rangle$, which lies at angle θ . Reflecting the vector at angle $-\theta$ about the line at angle θ produces a vector at angle $\theta + 2\theta = 3\theta$.

The composition of these two reflections is a rotation by 2θ . After k applications the state is at angle $(2k+1)\theta$ from $|B\rangle$, giving the stated formula. \square

The four diagrams below show the geometric effect of one Grover step. Angles are colour-coded: **blue** for θ between $|B\rangle$ and $|v\rangle$, **red** for θ between $|B\rangle$ and $Z_f |v\rangle$, and **orange** for the net 2θ rotation between $|v\rangle$ and $G |v\rangle$.

